

In-sensor optoelectronic computing using electrostatically doped silicon

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Abstract

Complementary metal-oxide-semiconductor (CMOS) image sensors are a visual outpost of many machines that interact with the world. While they presently separate image capture in front-end silicon photodiode arrays from image processing in digital back-ends, efforts to process images within the photodiode array itself are rapidly emerging, in hopes of minimizing the data transfer between sensing and computing, and the associated overhead in energy and bandwidth. Electrical modulation, or programming, of photocurrents is requisite for such in-sensor computing, which was indeed demonstrated with electrostatically doped, but non-silicon, photodiodes. CMOS image sensors are currently incapable of in-sensor computing, as their chemically doped photodiodes cannot produce electrically tunable photocurrents. Here we report in-sensor computing with an array of electrostatically doped silicon *p-i-n* photodiodes, which is amenable to seamless integration with the rest of the CMOS image sensor electronics. This silicon-based approach could more rapidly bring in-sensor computing to the real world due to its compatibility with the mainstream CMOS electronics industry. Our wafer-scale production of thousands of silicon photodiodes using standard fabrication emphasizes this compatibility. We then demonstrate in-sensor processing of optical images using a variety of convolutional filters electrically programmed into a 3×3 network of these photodiodes.

Main Text

Complementary metal oxide semiconductor (CMOS) image sensors have become an indispensable part of our data-driven world, where visual information prevails^{1,2}. The front-end silicon photodiode array in a CMOS image sensor converts light into electrical currents. These electrical data undergo analog-to-digital conversion and are then shuttled to a digital back-end for image processing. While this standard sequence of front-end image capture and back-end processing restricts the role of the photodiode array to sensing, emerging machine vision applications would benefit from data processing within the photodiode array itself^{3,4}. For example, in object tracking for self-driving vehicles, drones, or robots, where only the edges of objects are relevant^{5–8}, edge extraction in the front-end photodiode array would be much more economical in energy expenditure, processing latency, required bandwidth, and memory usage, as compared to transferring the whole image data containing superfluous information to the back-end digital processor—only to extract the edges⁹.

Such in-sensor computing would require an electrical modulation, or programming, of photocurrents. In fact, in-sensor computing has been recently demonstrated with electrostatically doped photodiodes whose photocurrents can be modulated with gate biasing^{10,11}. These pioneering works have realized electrostatically doped photodiodes by gating two-dimensional (2D) transition metal dichalcogenide (TMD) layers or their van der Waals (vdW) stacks^{12–14}. In contrast, such in-sensor computing is not possible with the present build of CMOS image sensors, for they employ chemically doped silicon photodiodes, whose photocurrents are not amenable to electrical modulation. Here, we report in-sensor computing with an array of electrostatically doped silicon *p-i-n* photodiodes, which can be seamlessly

integrated with the remainder of the CMOS image sensor electronics, while replacing the chemically doped silicon photodiode array. Such silicon-based approach could expedite the real-world application of in-sensor computing due to its compatibility with the mainstream CMOS electronics industry^{3,4,15,16}. Concretely, we first demonstrate large-scale device production by fabricating thousands of dual-gate silicon *p-i-n* photodiodes at the wafer scale. We then perform in-sensor computing on serial optical images using a 3×3 network of these electrostatically doped photodiodes by electrically programming the network into 7 different convolutional filters.

Electrostatically doped silicon photodiodes

The photocurrent of a diode, I_{ph} , grows with the power of the incident light, P , with the responsivity, R , being the proportionality constant¹⁷, *i.e.*, $I_{\text{ph}} = R \cdot P$. A conventional, chemically doped photodiode exhibits a constant responsivity R , since the parameters that determine R , especially the doping densities of the *p* and *n* regions, are fixed. On the other hand, in an electrostatically doped photodiode, where the doping densities can be modified by gate biasing, R is electrically programmable. The electrostatically doped photodiode can thus perform analog multiplication between the incident light power P and the electrically programmed responsivity R . This programmable optoelectronic analog multiplication is the key to in-sensor image processing.

Our electrostatically doped photodiode is built on an intrinsic silicon wafer. It contains two contact electrodes—*i.e.*, electrode 1 and 2—to provide the current path, and two top gate metals, which, when biased with the same voltage magnitude of opposite signs, create electrostatically doped *p* and *n* regions in silicon (Figs. 1a and b). The part of the silicon without any overlying gate metal is an intrinsic (*i*) region, and acts as a channel in the device. This channel region is directly exposed to light from above. The contact and gate electrodes are arranged in an interdigitated fashion for a high channel width/length ratio of 5576 μm /5 μm . Detailed fabrication steps are described in Methods and Supplementary Fig. 1. The resulting *p-i-n* diode exhibits a standard rectifying behavior (Supplementary Fig. 2), which confirms the electrostatic doping. As we swap the signs of the two gate biases, the rectifying behavior flips its polarity (Supplementary Fig. 2), which further verifies the electrostatic doping.

Illumination of the intrinsic channel region with a frequency of light higher than the silicon bandgap (~1.12 eV, or ~1,100 nm) generates a photocurrent. For this photocurrent generation mode, throughout this work, we bias both contact electrodes at zero voltage and define the current flow from electrode 1 to 2 as positive. The genesis of the photocurrent is the electrons and holes excited by the light, which are swept in opposite directions by the built-in potential (V_{bi}) of the diode, which is determined by the doping densities of the *p* and *n* regions (Fig. 1b). The electrostatic alteration of the doping densities via the gate voltages changes V_{bi} , which in turn can modify the magnitude and direction of the photocurrent for a given power of incident light. In other words, the gate biases tune the responsivity R .

We demonstrate this dependence of R on the gate biases by measuring the photocurrent with a fixed power, red-filtered halogen lamp that is periodically shuttered on and off, while the voltage at the gate above electrode 1 ($V_{G,1}$) is stepped up from -3 V to 3 V with a 0.5 V step, and simultaneously the voltage at the gate above electrode 2 ($V_{G,2}$) is stepped down from 3 V to -3 V with a 0.5 V step (inset of the Fig. 1c). The optical power of the light source, P_{source} , is 15 μW , which is different from, but proportional to, the power P of the light incident on the device, scaled according to the device and/or beam area. The measured photocurrent, shown in Fig. 1c, exhibits the expected modulation of R by the gate bias voltages. Repetition of such gate-controlled photocurrent modulation for ~ 50 min shows the stability of the programmability in R (Supplementary Fig. 3).

COMSOL Multiphysics simulation also confirms the operating principle of the electrostatically doped $p-i-n$ diode. The gating clearly creates p and n regions, with the band bending across the channel (Supplementary Fig. 4a-c) and the responsivity R changing with the gating as expected (Supplementary Fig. 4d; more on this shortly, in connection with Fig. 2c).

Programmable optoelectronic multiplication

We further investigate the dependence of the photo response on the gate voltages, and now, also on the light power (Fig. 2a). Figure 2b shows the photocurrent map with two gate voltages independently swept, each from -5 V to 5 V with a step of 0.1 V, while the photodiode is illuminated by blue laser light (473 nm) with a fixed P_{source} of 125 μW . When the two gate voltages are identical, *i.e.*, $V_{G,1} = V_{G,2}$, whether it is positive (n - n doping) or negative (p - p doping), no overall potential gradient develops, and thus no photocurrent should be produced. The corresponding pp to nn line, with zero current, is indeed close to the ideal positive diagonal line, and its slight deviation is possibly due to charge carrier trapping at defects formed during fabrication. On the other hand, when we sweep the two gate voltages at the same magnitude, but with opposite signs, along the negative diagonal line, the photocurrent monotonically increases from the negative maximum to the positive maximum, which is consistent with the monotonic change of V_{bi} from the negative maximum to the positive maximum (Fig. 1b). Figure 2c plots this photocurrent response along the negative diagonal line as a function of $V_{G,1} = -V_{G,2}$, which we denote as programming voltage V_p . This measured dependence of the photocurrent on V_p is also qualitatively consistent with the COMSOL simulation (Supplementary Fig. 4d). From here on, all the gate biasing is configured as $V_p = V_{G,1} = -V_{G,2}$.

Moreover, we demonstrate the linear dependence of I_{ph} on P_{source} —and therefore on P —for any given R programmed by tuning V_p . This linearity is important for high-fidelity analog multiplication between P and a given R . Figure 2d shows the measured I_{ph} as a function of P_{source} (red-filtered halogen lamp) for various V_p (and thus R) values. A simple linear fit yielding a high coefficient of determination (0.996 averaged across all V_p values) confirms the linear dependence of I_{ph} on P_{source} and thus on P , for

each programmed value of R . Linearity is also confirmed for different wavelengths of incident light (Supplementary Fig. 5).

Wafer-scale characterization of electrostatically doped silicon photodiodes

Electrostatically doped silicon photodiodes may accelerate the real-world realization of in-sensor computing due to their suitability for large-scale integration with CMOS electronics. To demonstrate, we have fabricated, in-house, 4,900 of the dual-gate p - i - n silicon photodiodes on a 4-inch silicon wafer (Fig. 3a, left) using the CMOS-compatible fabrication (see Methods). The fabricated wafer features $7 \times 7 = 49$ reticles, with each reticle containing $10 \times 10 = 100$ photodiodes (Fig. 3a, right).

Figure 3b shows photocurrent maps obtained by illuminating a 400 nm LED light with a fixed P_{source} of 170 μW serially—diode by diode—across an example reticle containing 100 photodiodes, for various V_p values (-5 V to 5 V with a 2 V step, clockwise from right, top corner). These maps show a high device-to-device uniformity in the responsivity programming within the reticle. In the wafer-scale photocurrent measurement of a 5×5 reticle array (2,500 photodiodes) with an automated probe station with V_p varied from -5 V to 5 V with a 0.1 V step, 2,372 devices showed programmable responsivity (~95% yield). Concretely, as we sweep V_p , the photocurrents of the 2,372 devices, in response to the 400 nm LED light with the fixed P_{source} of 170 μW , varied from -380 ± 50 nA to 430 ± 47 nA (Fig. 3c). Figure 3d shows the distribution of the 2,372 photocurrents for selected V_p values (-5 V to 5 V with a 1 V step), where device-to-device variations are more pronounced than those from the single reticle, which is standard at the wafer scale.

Optoelectronic convolutional image processing in a photodiode network

We connect 9 photodiodes as shown in Fig. 4a to perform analog multiplication between the incident light power and the programmed responsivity in each photodiode, and to sum, or accumulate, the resulting 9 photocurrents via Kirchhoff's current law. The photocurrent sum resulting from this analog multiply-accumulate (MAC) operation is a dot product between the 1×9 incident light power vector and the 1×9 vector of programmed responsivities. Consequently, the 9-photodiode network of Fig. 4a serves as an optoelectronic convolutional processor, with the 1×9 vector of programmed responsivities—or equivalently the 3×3 map of responsivities programmed across the photodiode array—serving as an image filter kernel. The accumulated photocurrent is converted to an output voltage (V_{out}) via a transimpedance amplifier on a printed circuit board (PCB). Our measurement system is detailed in Supplementary Fig. 6.

With an image filter kernel programmed, the 9-photodiode network not only captures an input scene, but also processes it simultaneously. Figures 4b-d show an example demonstration where the network finds

the edges of a moving light spot. We program the photodiode network to feature the specific responsivity map of Fig. 4b by independently tuning V_p of each photodiode (see Methods). This filter kernel is designed for edge-detection along the x -axis, resulting in positive and negative photocurrents when the photodiode network is at the right and left edges of the light spot, respectively, and otherwise negligible photocurrents. Figure 4c shows V_{out} monitored with a light spot from an LCD projector (power set to 255 out of 255, green channel only) moving from left to right at a frequency of 4 Hz, demonstrating the consistent positive (6 V) and negative (-6 V) responses as the spot moves over the array. We have evaluated this dynamic processing up to a spot movement frequency of 500 Hz (Fig. 4d), the maximum frequency of our optical setup.

Expanding from the simple example above, we perform in-sensor processing of a 256×256 pixel image (Fig. 5a, grey scale, 8 bit depth) with the contrast inversion filter kernel (Fig. 5b) programmed into the 9-photodiode network (see Methods for programming). Illumination of a 3×3 patch of the image onto the photodiode network using an LCD projector (green channel only) results in an accumulated photocurrent as the outcome of the optoelectronic convolution. By sliding the 3×3 patch through the 256×256 image and repeating the optoelectronic convolution, we generate a 254×254 matrix of accumulated photocurrents (a total of 64,516 accumulated photocurrents), which represents the image (Fig. 5c) processed with the contrast inversion filter kernel.

Besides the contrast inversion filtering, we have repeated in-sensor image processing using 6 other widely used filter kernels^{11,18–22}: difference of Gaussians (DoG), Gaussian blurring, image sharpening, box blurring, horizontal Sobel, and vertical Sobel filters (Supplementary Fig. 7). As the 63 photocurrent values programmed with a fixed P_{source} from the LCD projector (green channel only; 9 values per filter and a total of 7 filters), which correspond to the 63 programmed R values, are compared to their target values, which range from -2 μ A to 4 μ A, the maximum error was 18 nA. Since the ratio of the maximum error to the target range, 1/333, lies between $1/2^9$ and $1/2^8$, the programming accuracy is 8 bit. The images shined with the LCD projector (green channel only) and processed with these filter kernels are shown in Fig. 5e, bottom; the Sobel filtered image in Fig. 5e, bottom is a composite produced by the root sum of squares of the horizontal and vertical Sobel filtered images (Supplementary Fig. 8)²⁰. The juxtaposition of these images processed in the analog domain within the photodiode array (Fig. 5e, bottom) with those computed digitally (Fig. 5e, top) unequivocally verify our in-sensor computing scheme.

Conclusion

Bio-inspired computing has nucleated intense, worldwide research efforts in recent years, with in-memory computing motivated by the co-location of memory and computing in the brain, and with the even more recent in-sensor computing inspired by the sensory peripherals of the brain, where sensing is accompanied by early information processing^{3,4,10,11,23–28}. In this work, we have demonstrated analog image convolution processing as a form of in-sensor computing, where we have developed and utilized the electrostatically doped silicon photodiode array. This approach based on silicon devices suggests a

way to accelerate the practical realization of in-sensor computing by taking advantage of the mainstream CMOS electronics infrastructure. A monolithic integration of the electrostatically doped silicon photodiode array with conventional CMOS image sensor electronics, while replacing the chemically doped silicon photodiode array, is the next step of research our development suggests. The electrostatically doped photodiode we have fabricated in-house in this demonstration occupies an area orders of magnitude larger than a chemically doped photodiode in the state-of-the-art CMOS image sensor. Thus, increasing the density of the electrostatically doped photodiode array through substantial device miniaturization, while maintaining per-pixel gate and contact electrodes and their control by CMOS electronics, would be both a key direction and a challenge in this line of investigation.

Methods

Electrostatically doped silicon photodiode fabrication. Device fabrication began with wafer cleaning by dipping 4-inch un-doped Si wafer with 200 nm thermal oxide (Biotain Hong Kong Co., Limited, resistivity > 10,000 ohm·cm) in Piranha solution ($H_2SO_4:H_2O_2 = 3:1$) at 80 °C for 10 min. To build a device on a Si surface, the device fabrication area is defined through the conventional photolithography. Hexamethyldisiloxane (HMDS) is employed in advance to the photoresist to prevent undercut during a wet-etch process. Spin-Rinse-Dryer (SRD) is used for wafer drying through all the process. The 200 nm of thermal oxide is etched using buffered oxide etchant (BOE, 6:1) for 3 min. The pattern is stripped by acetone, isopropyl alcohol (IPA), and SRD. Contact electrodes (Cr/Au = 10/25 nm) are formed by conventional lift-off process using e-beam evaporator. Lift-off-resistor (LOR 3A) is employed for higher fabrication yield. 30 nm Al_2O_3 gate dielectric is deposited using atomic layer deposition (ALD) at 250 °C. After via formation by patterning with HMDS and BOE (1 min), gate electrodes (Cr/Au = 10/250 nm) are formed by the lift off process.

Iterative programming. While all the pixels are exposed to constant, maximum light from the LCD projector (255 for the 8 bit range, green channel only), V_p for each pixel is set to a calculated V_k for the k_{th} iteration calculated by the following equation:

$$V_k = (I_{target} - I_{k-2}) \times \frac{V_{k-1} - V_{k-2}}{I_{k-1} - I_{k-2}} + V_{k-2}, V_0 = 0 V, V_1 = 0.1 V$$

where I_{target} is the target current and I_k is the measured photocurrent at n_{th} iteration cycle, for each of 9 pixel. We keep modulating V_p until the difference between I_{target} and I_k is less than allowed error range, i.e., 23 nA if we set an 8 bit accuracy for the full range of 6 μA .

Declarations

Code availability

Experimental code is available from the corresponding authors on reasonable request.

Data availability

The data that support the findings of this study are available from the corresponding authors upon reasonable request.

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Author contributions

H.J., H.H., S.P., and D.H. conceived and designed the experiments. H.J., M.P., and S.L. designed and fabricated the electrostatically doped silicon photodiodes. H.H. designed the interface electronics. H.J., H.H., M.P. and S.L. performed the measurements of individual dual-gate *p-i-n* photodiodes. H.J. and H.H. performed the in-sensor image processing. M.L. and C.K. performed the wafer-scale measurements. W.J. performed COMSOL multiphysics simulations on the dual-gate *p-i-n* photodiode. H.J., H.H., W.J., M.L., C.K., M.P., S.L., and D.H. analyzed the data. S.P. and D.H. supervised the project. H.J., H.H., and D.H. wrote the manuscript. All authors discussed the results and implications, and reviewed the manuscript.

Competing interests

The authors declare no competing financial interests.

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Figures

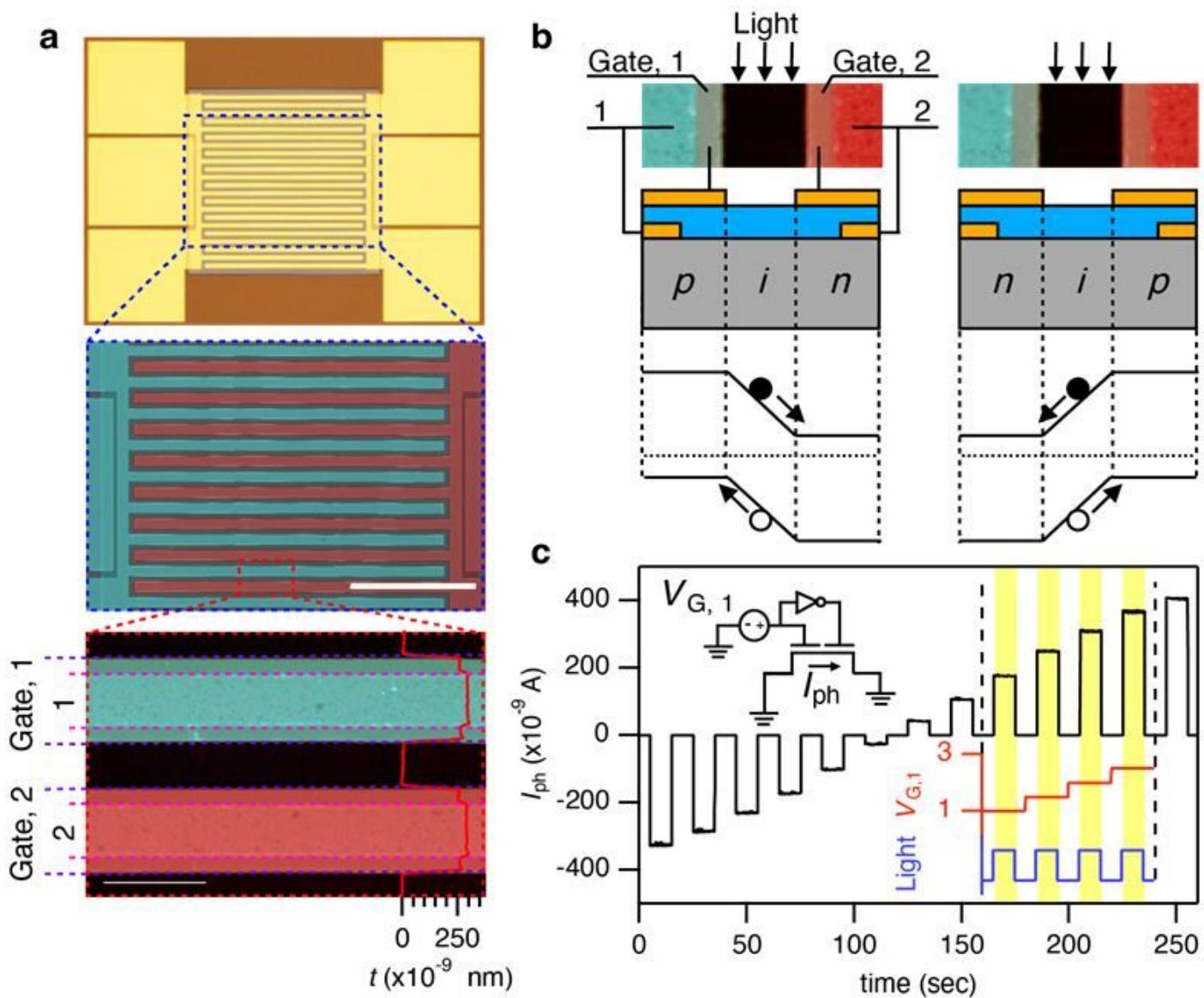


Figure 1

Electrostatically doped silicon *p*-*i*-*n* photodiode. **a**, Optical microscope (top), scanning electron microscope (SEM, middle), and atomic force microscope (AFM, bottom) images of an electrostatically doped *p*-*i*-*n* photodiode prototype. Contact electrode 1 and 2 and two gate electrodes above (false-colored with blue and red shades in the SEM and AFM images) are interdigitated. **b**, Part of the device SEM image (top) and corresponding schematic illustration (bottom) of the cross-sectional view of the photodiode, gate-biased to form *p*-*i*-*n* and *n*-*i*-*p* configurations. For a more realistic spatial profile of the electron concentration under gate biasing, see Supplementary Fig. 4. **c**, Measured photocurrent with pulsed light (blue) and stepped gate voltages ($V_{G,1}$, red; $V_{G,2}$, not shown). A red-filtered halogen lamp ($P_{source} = 15 \mu\text{W}$) is used as the light source. $V_{G,1}$ is stepped up from -3 V to 3 V with a 0.5 V step, and $V_{G,2}$ is simultaneously stepped down from 3 V to -3 V with a 0.5 V step.

Figure 2

Programmable photo response of the dual-gate silicon *p-i-n* photodiode. **a**, Schematic illustration of the measurement setup. Incident light with power P —*i.e.* P_{source} scaled according to the device and/or beam area—is converted to the photocurrent, I_{ph} , which is modulated by the two gate voltages, $V_{G,1}$ and $V_{G,2}$. **b**, Photocurrent map measured with each gate voltage independently swept from -5 V to 5 V with a 0.1 V step. **c**, Photocurrent response with $V_p = V_{G,1} = -V_{G,2}$ swept from -5 V to 5 V with a 0.1 V step. The light source for parts **b** and **c** is a blue laser (473 nm) with a P_{source} of 125 μW . **d**, The measured photocurrent vs. P_{source} with V_p as a parameter, varied from -4 V (blue) to 4 V (red) with a 1 V step. The light source is a red-filtered halogen lamp.

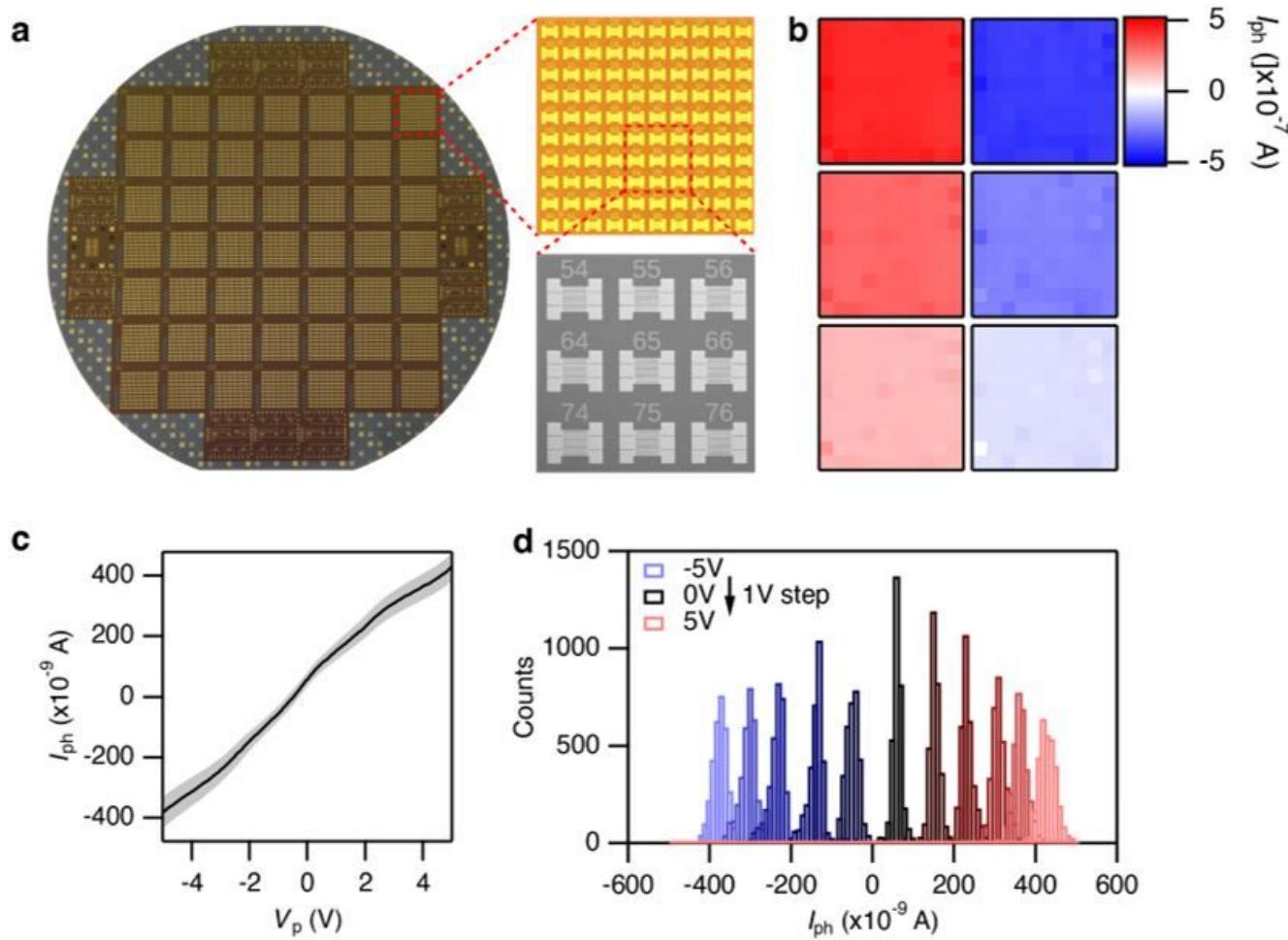


Figure 3

Wafer-scale array of the dual-gate silicon *p-i-n* photodiodes. **a**, Optical images of a fabricated wafer containing 7×7 reticles (left) and a reticle containing 10×10 photodiodes (top, right), and an SEM image of 9 example photodiodes. **b**, The photocurrent map of a single reticle (10×10 photodiodes) with V_p varied from -5 V to 5 V with a 2 V step (clockwise from top, right). **c**, The average (solid line) and standard deviation (shades) of photocurrents measured from 2,372 working photodiodes from 5×5 reticles containing 2,500 photodiodes, with V_p varied from -5 V to 5 V with a step of 0.1 V. **d**, The histogram of the photocurrent data collected in part **c**, with V_p from -5 V (blue) to 5 V (red) shown in 1 V increments for clarity. All measurements shown in this figure were performed with a 400 nm LED light with a P_{source} of 170 μW .

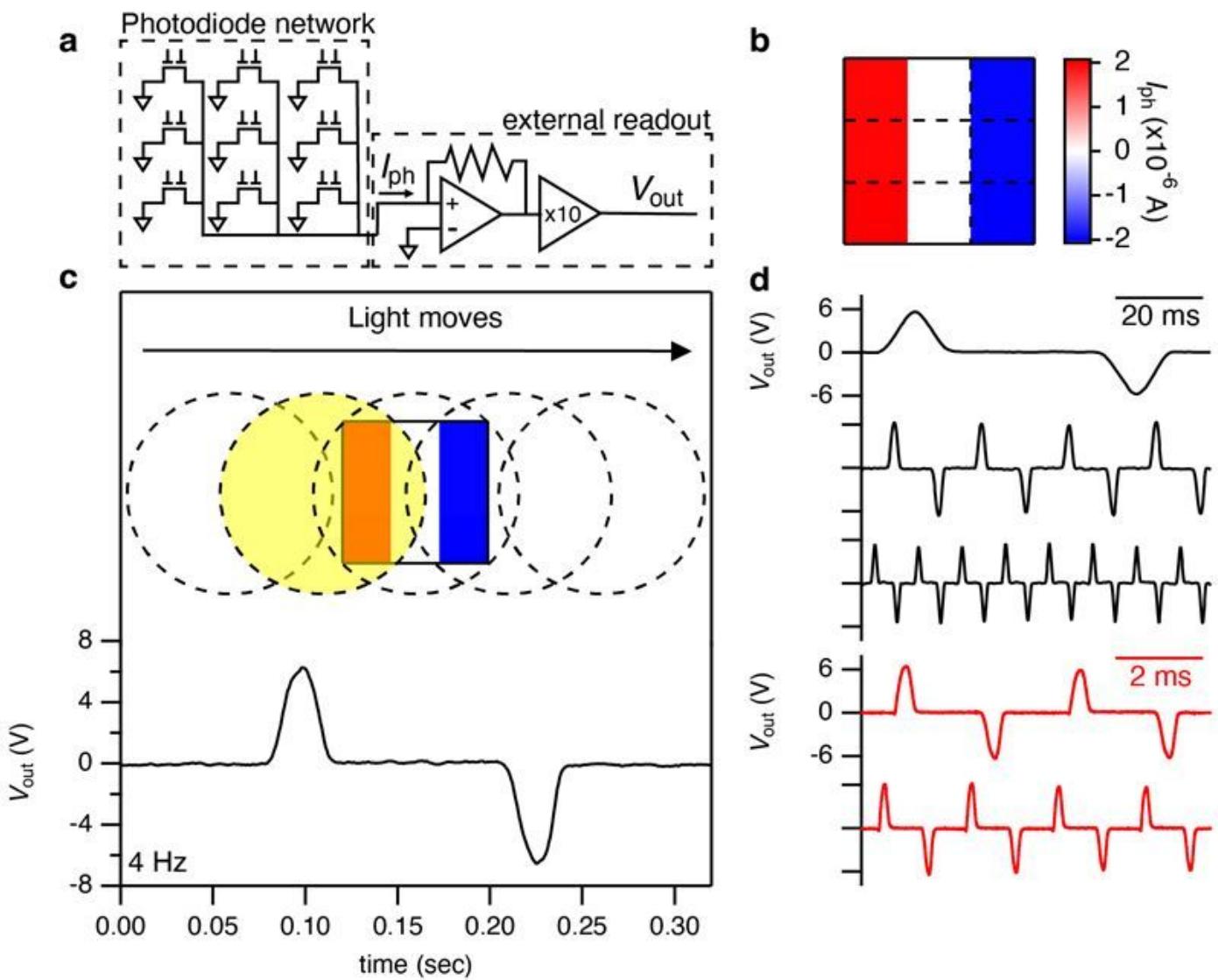
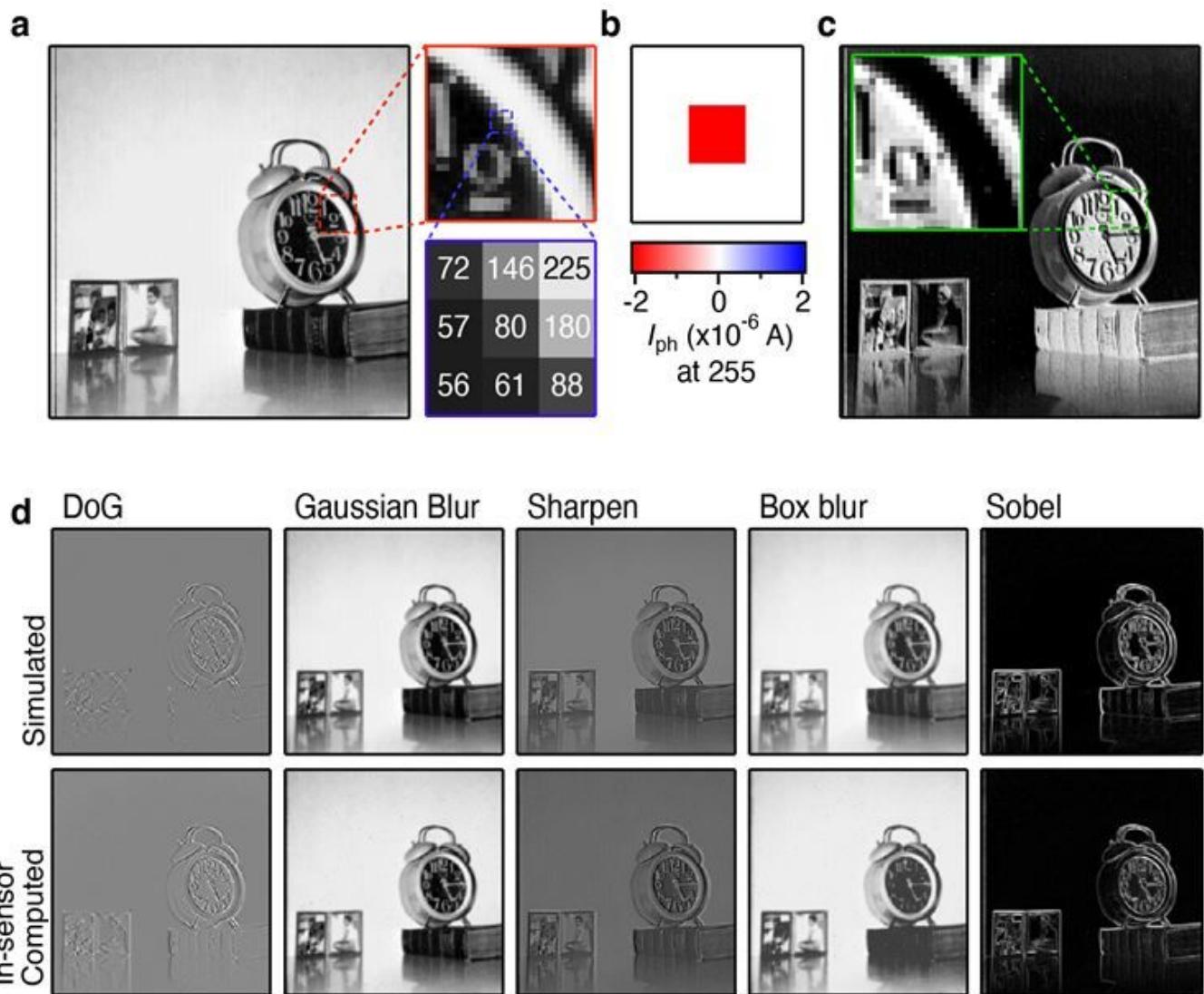


Figure 4

A 3×3 photodiode network for analog multiply-accumulate (MAC) computation. **a**, Schematic illustration of the network comprising 9 dual-gate silicon *p-i-n* photodiodes. The accumulated photocurrent as a result of the analog optoelectronic MAC operation is converted to a voltage (V_{out}) by a transimpedance amplifier on a printed circuit board. **b**, A photocurrent map programmed with a constant light power, *i.e.*, a responsivity map, which represents a filter kernel for edge-detection along the x -axis. **c**, Measured V_{out} of the photodiode network arranged into the filter kernel of part **b**, with a light spot moving from left to right at 4 Hz. **d**, Repetition of part **c**, but with the light spot moving frequencies at 10, 50, 100, 250, and 500 Hz. For all experiments in this figure, an LCD projector (green channel only) with a power 255 out of 255 is used as a light source.

**Figure 5**

In-sensor image processing using the 3×3 dual-gate $p-i-n$ photodiode network. **a**, A 256×256 input image (left) and its example portion (top, right). The bottom right is an example of a 3×3 patch from this input image, which is projected onto the photodiode network. **b**, A programmed photocurrent map with a fixed power of light—*i.e.*, a responsivity map—for contrast inversion filtering. The maximum LCD projector brightness (255 out of 255, green channel only) is used for this programming. **c**, The 254×254 map of accumulated photocurrents with the 9 photodiode network programmed as in part **b**, where the 64,516 accumulated photocurrents are serially obtained by illuminating, using the LCD projector (green channel only), the photodiode network with a 3×3 patch sliding through the 256×256 input image. **d**, Various filtered images obtained with digital computing (top) and in-sensor computing (bottom).

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