On the use of graphene to improve the performance of concentrator III-V multijunction solar cells

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Abstract
Graphene has been intensively studied in photovoltaics focusing on emerging solar cells based on thin films, dye-sensitized solar cells, quantum dots, nanowires, and so forth. However, the typical efficiency of these solar cells incorporating graphene is below 16%. Therefore, the photovoltaic potential of graphene has not yet been shown. In this work, the use of graphene for concentration applications on III-V multijunction solar cells, which indeed are the solar cells with the highest efficiency, is demonstrated. First, a wide optoelectronic characterization of graphene layers is carried out. Then, the graphene layer is incorporated onto triple-junction solar cells, which decreases their series resistance by 35% (relative), leading to an increase in fill factor of 4% (absolute) at concentrations of 1000 suns. Simultaneously, the optical absorption of graphene produces a relative short-circuit current density decrease in the range of 0% to 1.8%. As a result, an absolute efficiency improvement close to 1% at concentrations of 1000 suns was achieved with respect to triple-junction solar cells without graphene. The impact of incorporating one and two graphene monolayers is also evaluated.

KEYWORDS
concentrator photovoltaics, graphene, III-V solar cells
1 | INTRODUCTION

Aiming at the reduction of the cost of concentrator photovoltaic (CPV) systems, the use of ultrahigh concentration levels (≥1000 suns) in III-V multijunction solar cells (MJSCs) has been proposed.\(^1\) Although very high-efficiency MJSC designs are being developed for this purpose (with current champion efficiencies around 44%-46%),\(^4\) their series resistance is responsible for important efficiency losses when operating under highly concentrated light.\(^7\) The most important components of the series resistance in a MJSC are those of the front metal grid and top cell emitter/window, given the fact that tunnel junctions have demonstrated efficient operation up to several thousand suns.\(^6\) Moreover, when MJSCs operate inside an optical concentrator, spatial non-uniformity of the light impinging on the solar cell aggravates series resistance losses.\(^9\) To overcome these problems, the use of graphene in concentrator III-V MJSCs is proposed in this work. Graphene is incorporated covering both the front grid metallization and the topmost semiconductor layer to improve photocurrent extraction and therefore, the solar cell efficiency.

Since the first isolation of a single layer of graphene in 2004,\(^10\) its unparalleled properties have allowed its integration in a wide variety of electro-optical applications (LEDs, solar cells, advanced sensors, etc.).\(^11\) In particular, in the field of solar cells, the use of graphene has been intensively studied in emerging architectures based on thin films, dye-sensitized solar cells, quantum dots, nanowires, perovskites, and so forth, resulting in efficiencies typically ranging 0.26%-16.2%.\(^14\) However, these efficiency figures are strongly affected by inner limitations in the efficiency of these solar cells. Therefore, the real potential of graphene (as an advanced energy material) to achieve very highly efficient solar cells is still to be evinced.

Accordingly, in this paper, we face the use of graphene in one of the most challenging types of solar cells, namely, concentrator III-V multijunction solar cells. The defying of these solar cells derives from the following facts: a) this type of cells exhibits the highest power conversion efficiency,\(^4\) b) they are currently fabricated by means of a mature and robust technology at industrial level,\(^7\) and c) their operation at concentration levels of 1000 suns and beyond\(^21\) is an excellent benchmark to demonstrate the use of graphene to improve the photocurrent extraction and thus, to increase the efficiency of concentrator III-V MJSCs. Although the experimental incorporation of graphene in this work has been implemented on lattice matched GaInP/Ga(In)As/Ge triple-junction solar cells (3JSCs), the use of graphene we are proposing might be applied to the wide palette of III-V MJSC architectures available,\(^4,6,22,23\) as well as to other types of solar cells.

Through the paper, we first analyze the specific requisites and properties that graphene must fulfill for its successful integration into concentrator MJSCs. Then, we characterize the key optoelectronic properties of graphene layers to ensure their suitability in concentrator MJSCs. After that, we describe the experimental procedure to manufacture MJSCs with graphene. Finally, we compare the experimental performance of 3JSCs with and without graphene, which results in almost 1% (absolute) increase in the power conversion efficiency at around 1000 suns which increases as concentration rises. This represents a lower bound for the efficiency gain, therefore, future improvements that could be added to concentrator III-V MJSCs exhibiting efficiencies around 40% are discussed.

1.1 | Requisites for graphene in concentrator III-V MJSCs

In this paper, we use graphene on top of concentrator GaInP/Ga (In)As/Ge 3JSCs in conjunction with a conventional solar cell front metal grid (see Figure 1). The goal is that graphene creates paths of low electrical resistance for the photocurrent extraction in parallel to those of the emitter, window, and metal grid. These additional paths should result in a decrease of the series resistance which would, in turn, increase the fill factor (FF) and open-circuit voltage ($V_{oc}$). Simultaneously, the optical transmission of graphene needs to be very high since otherwise the efficiency improvements achieved by the decrease in series resistance would be counterbalanced by the optical losses which would reduce the short-circuit current density ($J_{sc}$).

Among the series resistance components of the 3JSC without graphene, the most important ones are schematized in Figure 1B, which include the following:

- the sheet resistance of the front metal grid ($R_{fg}$) with typical values of 0.05-0.20 $\Omega$/□, which are good enough for concentrator solar cells with a typical size in the unit of mm\(^2\) required for a reduction, among others, of the series resistance and a low overheating at very high light concentrations\(^1\);
- the specific contact resistance between the front metal grid and the n'-GaAs cap layer ($R_{fc}$) with values ranging $10^{-5}$-$10^{-6}$ $\Omega$-cm\(^2\);
- the top cell emitter sheet resistance ($R_{e, top sheet}$) with values of 300-500 $\Omega$/□ (which includes both the AlInP window ($R_{w}$) and emitter sheet resistance ($R_{e}$)) derived from a trade-off of reducing the series resistance while keeping a high photogeneration.

When the same 3JSC incorporates graphene, three new paths for photocurrent extraction might participate, namely, $R_{fg}$, $R_{gw}$, and $R_{gm}$ (encircled in red in Figure 1B). The requisites for these resistances are the following: a) low sheet resistance of graphene to reduce the value of the lateral graphene resistance, $R_{lg}$; b) low contact resistance between graphene and the front metal grid to minimize $R_{gm}$, and c) low contact resistance between graphene and the top semiconductor (window layer) to get a low enough $R_{gw}$.

The graphene sheet resistance, $R_{g}$, has to be as low as possible. In order to obtain the maximum benefit from the incorporation of graphene, a reasonable initial target is to aim for values equal or lower than the emitter sheet resistance, $R_{e, top sheet}$, that is, $300$-$500$ $\Omega$/□. These values are achievable by Chemical Vapor Deposition (CVD)-grown graphene.\(^24\) We have measured (see the Experimental section) a moderate mobility in the range of 750 to 2000 cm\(^2\)/V·s in our graphene layers and a resistivity in the range of 2.4 · $10^{-5}$ to 5.7 · $10^{-6}$ $\Omega$-cm, resulting in a sheet resistance ranging from 200 to 500 $\Omega$/□, which complies with the requirement of being in
the range of $R_{t}$ field. Further reduction in sheet resistance can be achieved by chemically doping graphene.\textsuperscript{25-27} In relation to the graphene-front grid metallization interface, the requirement is to achieve an ohmic contact resulting in $R_{g-m}$ values around the specific front contact resistance between the front metal grid and the n$^+$-GaAs cap layer ($R_{f}$). Figure 2A shows the test structures

**FIGURE 1** Graphene integration into a concentrator III-V MJSC. Graphene is deposited in contact with both the front grid metallization and the semiconductor active area. A, Graphene (black hexagons) on top of both the metallization (fingers + busbar, in yellow) and semiconductor active area (AlInP window layer highlighted in cyan). B, Zoomed view of a small portion of panel a, where all the resistances participating in the front extraction of photocurrent have been included. Graphene integration provides alternative routes for current extraction marked as resistances encircled in red which represent the following: 1) the contact resistance between graphene and window layer ($R_{g-w}$); 2) the lateral graphene resistance ($R_{g}$), and 3) the contact resistance between graphene and metal grid ($R_{g-m}$). [Colour figure can be viewed at wileyonlineibrary.com]

**FIGURE 2** Characterization of the graphene-metal contact. A, Test structures used for the measurement of the contact resistance between graphene and the front metallization made of AuGe/Ni/Au and Ti/Pd/Au. Black arrows indicate the contacts used for current injection and voltage measurement. Configuration A_01 measures the I-V curve through AuGe/Ni/Au-graphene-Ti/Pd/Au, while Configuration A_02 measures the I-V curve through Ti/Pd/Au-graphene-Ti/Pd/Au. B, Experimental I-V curves for both Configuration A_01 and Configuration A_02. [Colour figure can be viewed at wileyonlineibrary.com]
manufactured to check the graphene-metal contact (in our 3JSCs the front metal grid is made of AuGe/Ni/Au), while Figure 2B shows an example of the I-V curves measured, which confirm an ohmic behavior. The results of the TLM (transmission line method) analysis (detailed in the Experimental section) yield a specific contact resistance of the graphene-metal contact of 2 · 10⁻⁴ Ω cm², which is close to values reported in the literature.²⁸

Regarding the equivalent resistance of the graphene-AlnP window layer interface, \( R_{g-w} \), we firstly simulated its band alignment (Figure 3B), which showed a high dependence on the AlnP doping level. Subsequently, we calculated the dark J-V curve of a 3JSC with an n-AlnP window layer of 40 nm with different doping levels and covered by a graphene layer with different work functions. Figure 3B shows the simulated dark J-V curves taking into account the influence of both thermionic emission and tunneling mechanisms. Irrespective of AlnP doping level and for the three graphene work functions considered, tunneling appears as the dominant mechanism; its influence rises as the doping level increases. Figure 3B shows that the impact of the graphene work function is very dependent upon the AlnP doping level. In fact, if the AlnP doping level increases over 10¹⁹ cm⁻³, dark J-V curves are largely insensitive to variations in the graphene work function (which is affected by surface condition, that is, contamination or oxidation of the AlnP window layer may vary the work function value) for the whole voltage range studied (i.e., the dashed black line, solid black line and squares virtually overlap in Figure 3B). In such situation, current transport through the graphene-AlnP interface resembles an ohmic contact. Since the experimental doping level of the AlnP window layer of the 3JSCs in this study is in the range of 2 · 10¹⁹ to 3 · 10¹⁹ cm⁻³, we may consider that our graphene-AlnP contact (assumed as ideal, i.e., without any interface contamination such as oxides) will behave as ohmic.

Further reductions on the series resistance are possible by stacking several graphene layers on top of each other. Ideally, the larger the number of graphene layers, the lower the lateral series resistance.

![Image](image_url)

**FIGURE 3**  Electrical and optical performance of the graphene-AlnP interface. A, Simulated energy band alignment at equilibrium for a graphene monolayer (with a work function of 4.86 eV, which was experimentally measured) and an AlnP window layer as function of its doping level. B, Simulated dark J-V curves of a 3JSC integrating one graphene monolayer as a function of the graphene work function and the AlnP doping level. Simulations consider three values for the graphene work function to take into account the experimental values of our graphene measured by UPS, namely, 4.86 eV (dashed) and 4.95 eV (solid) and also the ideal value of 3.78 eV (squares) coincident with the electronic affinity of AlnP. Different equivalent light concentration levels (1-1000 suns) are included as guides to the eye by considering a \( J_{sc} \sim 14 \text{ mA/cm}^2 \) at 1 sun. C, Experimental transmittance (T) of AlnP samples without graphene and with 1GML and 2GML. D, Relative transmittance of 1GML and 2 GML deposited on AlnP samples (Equation 1). The uncertainty in the optical transmittance has been assumed to be equal to the standard deviation [Colour figure can be viewed at wileyonlinelibrary.com]
but the lower the optical transmission too. In this way, a trade-off arises around the optimum number of graphene layers, which will be discussed below.

Regarding the optical behavior, we have studied the optical impact of transferring graphene on AlInP window layers by measuring the spectral reflectance (R) and transmittance (T), as indicated in the Experimental section. In this work, we have used graphene grown by CVD, as this growth technique offers excellent control of the material properties over large areas and it is scalable to industrial levels.

Figure 3C shows the experimental transmittance of the AlInP samples without graphene and with one and two graphene monolayers (1GML and 2GML). The simulated transmittance curves using the generalized transfer matrix method with the refractive index and extinction coefficients determined for our CVD graphene (not shown in Figure 3C for the sake of clarity), fit perfectly the experimental ones for the AlInP structures with and without graphene. As expected, transmittance decreases as graphene thickness increases for all wavelengths.

The relative transmission of graphene layers ($T_{R,GML}(\lambda)$) can be obtained using the measured transmittance of the AlInP sample ($T_{AlInP}(\lambda)$) together with the transmittance of the AlInP samples covered with 1GML and 2GML ($T_{GML+AlInP}(\lambda)$):

$$T_{R,GML}(\lambda) = \frac{T_{GML+AlInP}(\lambda)}{T_{AlInP}(\lambda)}$$ (1)

Figure 3D depicts these results. The resulting average transmittance for 1GML and 2GML within the 400-1700 nm range (which is of interest for the kind of 3JSCs of this work) is, 98.7% and 98.0%, respectively. From the transmittance curves of Figure 3D and the corresponding absorbance for 1GML and 2GML whose average values are 1.2% and 3.7%, respectively for the visible range (400-700 nm). These values are consistent with those of Lee et al. who predicted by band theory an absorbance of 1% for graphene on silicon substrates, which have a refractive index of 3.4, close to that of AlInP. It should be noticed that this is a different case from the 2.3% absorbance of suspended graphene.

### 1.2.2 Graphene wet transfer for its characterization

Graphene wet transfer was also carried out following the procedure from Kong’s group. In the wet transfer method, a polymer film is deposited on the graphene/Cu stack and then the sample is baked at 130°C for 10 min. Immediately after that, the catalyst metal is etched away with a Cu etchant. The resulting PMMA/graphene stack is cleaned in HCl:H2O. These steps help remove some debris from the Cu etchant such as oxidized metal particles. The PMMA/graphene film is then scraped, deposited on the substrate of interest, and dried with N2.

Afterward, baking steps are carried out to promote graphene adhesion on the semiconductor and then PMMA is removed by soaking the sample in acetone. This step sequence does not ensure a completely clean surface since, in the end, graphene has been exposed to 1) the metal etchant, 2) polymer residues, and 3) water and dust particles that could adhere during the different transferring steps and eventually strongly affect the final properties of the material. Therefore, a final thermal treatment is used to get rid of such contaminants.

### 1.2.3 Electro-optical characterization of CVD graphene

The transmission line method (TLM) and Van der Pauw measurements were carried out on CVD graphene monolayers for their electrical characterization (graphene sheet resistance, mobility, and doping levels among others). Graphene monolayers were transferred by the wet transfer method into SiO2/Si substrates. TLM patterns were fabricated onto 1GML/SiO2/Si using e-gun evaporation of 5.5 nm Ti/45 nm Pd/110 nm Au. Conventional photolithographic techniques were used to define pads with an area of 400 μm × 150 μm and with pad distances of 90/180/360/720/1440 μm. In order to get rid of the lateral deviation of the current, a second photolithography step for TLM isolation was needed. A dry etching was applied to etch graphene from the desired areas.

Van der Pauw measurements were performed after the graphene transfer to SiO2/Si substrates using a homemade set-up. For simplicity, no photolithography was applied. Instead, a silver paint was used to form an ohmic contact pad by applying it on the four edges of a square-shaped graphene layer. By applying a current of 1 mA and a magnetic field of 2900 Gauss, the Van der Pauw technique results in the determination of the graphene sheet resistance, mobility, and carrier density.

Graphene work function in contact with the AlInP window layer was obtained using ultraviolet photoelectron spectroscopy (UPS). A
photon source of He (I), 21.21 eV, was used. Since work function is affected by the surface condition, that is, contamination or oxidation may have a significant influence on the work function value obtained, a previous Argon (Ar) cleaning treatment was performed for some samples to see that influence. This cleaning treatment consisted in an Ar⁺ sputtering with an energy of ~1 keV for 60 s. Graphene/AlnP work function with and without cleaning treatment resulted in 4.95 eV and 4.86 eV, respectively.

Spectral transmittance of the topmost part of the 3JSCs cells was determined by measuring the stack consisted of graphene (1GML and 2GML) transferred onto an AlnP window layer stack with a transparent epoxy on top of a glass holder. Prior to graphene transfer, transmittance of AlnP window layer on the glass holder was measured. Therefore, the relative transmittance of graphene ($T_{R,GML}(\lambda)$) was obtained using the transmittance of the AlnP window layer on top of a glass holder for its normalization following Equation (1). Spectral transmittance measurements from visible- to middle-infrared range were done by means of a Perkin Elmer lambda 1050 Spectrometer.

1.2.5 | Integration of graphene on concentrator 3JSCs

Graphene monolayers synthesized by CVD were transferred by the already-described wet transfer method on triple-junction solar cells. In order to achieve a fair comparison between 3JSCs with and without graphene, CVD-graphene was deposited only on one-half of the epiwafer so the epistucture and device processing were identical for all solar cells with and without graphene. Additionally, when a second monolayer (2GML) of graphene was deposited, a subsequent wet transfer process was repeated. Although thermal treatments between 300 and 350°C are desired as the last step of the graphene transfer process, during this work, the presence of both front and back metalization on the solar cell wafer limited this temperature down to 200°C to avoid their deterioration.

1.2.6 | Singulation of 3JSCs incorporating graphene (RIE + mesa etching + dicing)

One of the last steps of the solar cell manufacturing is the electrical isolation of each individual device in the wafer by mesa etching. Since graphene is chemically inert to wet etchants and is placed covering the semiconductor layers (see Figure 4), a purely chemical mesa etching will not be effective if graphene is not previously removed from the mesa trenches. Thus, prior to mesa isolation, a second photolithography step was required to protect graphene at areas where it has to be preserved (i.e., on the active area and on the front grid metallization). Graphene was etched away from the isolation trenches via RIE and afterward, the mesa etching was carried out. Finally, the epiwafers containing the solar cells were rinsed in acetone and methanol to remove the photoresist. After that, solar cells with and without graphene were diced and individually packaged onto printed circuit boards (PCBs). No ARC was deposited on the devices to facilitate the comparison between solar cells with and without graphene as well as to determine more clearly the graphene optical impact on 3JSCs. Figure 4 visually summarizes the key steps in the manufacturing of 3JSCs with graphene.

1.2.7 | Characterization of graphene-based 3JSCs

The characterization of 3JSCs with and without graphene consisted of the dark I-V, external quantum efficiency (EQE), concentrator response, and electroluminescence (EL) mapping. I-V curves were measured with the four-point probe technique to suppress the influence of the series resistance of the probe and wires on the measurements. The set-up employed for the EQE consists of a Xe lamp used as white light source, which passes through a Horiba Jobin Yvon monochromator (TRIAX180) and a filter wheel. Further details on the system and the measurement can be found in. Measurements under concentration were performed with a flash simulator. For a detailed explanation of the concentration set-up, the reader is referred to. Regarding EL measurement, driving the cells into forward bias makes them behave as light emitting diodes (LEDs). A current injection of 100 mA was used to emulate the solar cell current photogeneration at concentration levels around 900x, (assuming $J_{SC} > 10^{-11}$ mA/cm² without ARC). The emitted light is recorded with a CCD camera (CCD4 MV1-D1312I Photon Focus) in a home-made EL set-up. The electroluminescence images investigated in this work were those of the GaInP top subcell only (emission at 650-690 nm), for which a short pass filter is used (cut-off wavelength 800 nm).
2 | RESULTS AND DISCUSSION

In order to characterize the solar cells with and without graphene, we firstly checked that graphene layer conformally follows the solar cell surface by means of atomic force microscopy (AFM) in tapping mode. As Figure 5A left shows, the adherence of graphene between metal fingers and the solar cell active region is reasonable and the only regions where the graphene layer does not touch the active region is next to the metal fingers. In the case of the busbar (Figure 5A right), the abrupt step denotes a better adherence between graphene and active area. Therefore, we can expect that a significant part of the solar cell surface is in an intimate contact with the graphene layer. This assumption is confirmed by EL map measurements of 3JSCs without graphene and with 1GML (Figure 5B) which show similar performance indicating that the solar cell with 1GML has no regions with poor contact between the graphene layer and the solar cell surface. A more careful examination of Figure 5B shows that the cell with 1GML exhibits at the center a higher EL emission (less cyan color) than the cell without graphene (as Figure 5C highlights), which anticipates a better current distribution thanks to graphene.

The EQE is presented in Figure 6A for the GaInP top, Ga (In)As middle and Ge bottom subcells of the 3JSCs with and without graphene. As it can be observed in Figure 6A, a very small decrease in EQE for some wavelengths is detected after graphene integration. To quantify such decrease, the short-circuit current density ($J_{sc}$) was calculated for several solar cells from the convolution of their EQE with the reference spectrum for concentration (AM1.5d ASTM G173-03). That results in a relative $J_{sc}$ decrease with respect to solar cells without graphene in the 0%–1.8% range for 1GML while 2.0%–3.8% for 2GML, which are consistent with the average absorbance of 1.2% and 3.7% for 1GML and 2GML, respectively, described above. Concentration measurements at different irradiance levels were performed on the same solar cells without ARC in order to ease the
comparison between solar cells with and without graphene. Figure 6. B-D shows the evolution of FF, V_{oc}, and power conversion efficiency (\eta) as a function of concentration. Figure 6B exhibits FF values similar for solar cells with and without graphene at low concentrations, thus, the influence of graphene is almost negligible. It is not until medium concentrations (\geq 300\times) when FF of solar cells with graphene is higher. In fact, FF improvement increases with concentration, reaching an increment of around 4% (absolute) at approximately 1000\times. This can be explained by the decrease in the series resistance caused by graphene, which minimizes the ohmic losses during the photocurrent extraction, with the subsequent improvement in FF. The FF data fitting to our solar cell models show a series resistance of 34 \Omega\cdot\text{cm}^2 for the 3JSCs without graphene which is reduced to 22 m\Omega\cdot\text{cm}^2 for the cells with 1GML (i.e., a 35% relative reduction of series resistance). The
addition of a second graphene monolayer results in an additional FF improvement. Nevertheless, this second improvement is lower than that expected for 1GML + 1GML. A possible explanation for this fact could be ascribed to the effect of contaminants (PMMA, water molecules, etc.) from the wet transfer process that might have been trapped between the two stacked graphene monolayers (as was reported in previous works) hindering photocurrent transfer between them.

Figure 6C shows the evolution of Voc versus concentration for the same 3JSCs with 1GML, 2GML, and without graphene showing a very similar performance. Figure 6d shows that efficiency of 3JSCs without graphene is higher at low concentrations but is outdone at high irradiance levels by 3JSCs with graphene. The difference at high concentrations shows up clearly in the I-V curves of Figure 6E. This is because FF and Voc improve with the use of graphene thanks to the series resistance reduction, but at the same time, Jsc slightly decreases because of graphene optical absorption. This trade-off leans toward FF and Voc only for high enough concentrations. As an example, at 1000×, absolute efficiency gains close to 1% (1GML) and 0.5% (2GML) are achieved, with respect to 3JSCs without graphene and these efficiency gains are even higher at concentrations beyond 1000 suns. This fact can be explained in terms of the transparency-resistance trade-off: for 2GML the decrease in resistance obtained in our experiment does not improve the FF so much as to counterbalance the higher optical losses. Of course, this could change if better photocurrent transfer would have been achieved between the two graphene monolayers but until this problem is not experimentally solved, our work shows that the optimum number of graphene layers is one. It is worth mentioning that none of the 3JSCs have ARC so, assuming a typical power conversion efficiency increase of 35% due to ARC, these solar cells with ARC would result in absolute efficiencies higher than 39% at 1000 suns. Each point of all the figures is the average of three different solar cells. (a) EQE of the GaInP, Ga(In)As, and Ge subcells of a 3JSC without graphene and with 1GML and 2MLG. The Ge subcell with 2MLG is not shown. b-e, Performance as a function of light concentration of 3JSCs without and with graphene (with 1GML and 2GML). (b), Fill factor (FF). (c), Open-circuit voltage (Voc). (d), Power conversion efficiency (η). (e), I-V curves at the highest concentration in each 3JSC design. Solar cells have no ARC [Colour figure can be viewed at wileyonlinelibrary.com]
these solar cells with ARC would result in absolute efficiencies higher than 39% at 1000 suns.

To further improve the results of this paper, the deposition of graphene on top of the window layer could be done prior to the evaporation of the front grid metallization. In this way, graphene will be deposited directly on a completely planar and clean surface without the metal fingers. This would reduce both the resistance of the graphene-AllnP window layer interface, $R_{g-w}$, and the resistance of the graphene-front grid metallization interface, resulting in a higher FF. Further reductions in graphene sheet resistance can be achieved by chemical doping. Besides, this alternative approach of integrating graphene prior to the front grid metallization would allow a higher flexibility in the choice of the annealing time and temperature of graphene for a more effective suppression of surface residue. This would result in a reduction of the optical absorption of graphene with the subsequent increase of $J_{sc}$. Therefore, these increases of FF and $J_{sc}$ will produce an improvement in the power conversion efficiency of concentrator III-V MJSCs.

3 | SUMMARY AND CONCLUSION

We have assessed CVD-graphene incorporation atop concentrator III-V MJSCs. First, the requisites and properties for graphene to be integrated in those cells were stated and then a wide optoelectronic characterization of graphene layers was carried out in order to confirm that our CVD graphene complies with such requirements.

High-concentrator GaInP/Ga (In)As/Ge triple-junction solar cells were manufactured on epifaces grown by MOVPE. Graphene monolayers were transferred onto one half of the 3JSC structures using the wet transfer method. In this way, a fair comparison between 3JSCs with and without graphene can be done because their epistructures and device processing were identical. No ARC was deposited in order to ease that comparison.

Graphene incorporation into triple-junction solar cells shows an effective decrease of the series resistance (35% relative), leading to an increase in FF of 4% (absolute) at concentrations of 1000×. EL measurements confirm a better spreading of the photocurrent in solar cells with graphene. Simultaneously, the optical absorption of graphene produces a relative $J_{sc}$ decrease in the range of 0% to 1.8%. As a result, an absolute efficiency improvement close to 1% at concentrations of 1000× was achieved with respect to 3JSCs without graphene and this improvement increases with concentration. The impact of one and two graphene monolayers was also evaluated showing that a single graphene layer is better than two. These results highlight the potential of graphene to reduce the impact of series resistance in concentrator MJSCs and pave the way to the exploitation of CPV designs working over a thousand suns.

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